



Complex valued waves are separable:

$$l(\mathbf{y}) = \exp(i(\boldsymbol{\omega}^T \mathbf{y}))$$

= $\exp(i\omega_1 y_1 + i\omega_2 y_2))$
= $\exp(i\omega_1 y_1) \cdot \exp(i\omega_2 y_2)$
=: $l_1(y_1) \cdot l_2(y_2)$

- The same is not true for real valued waves.
- Separability can be used to prove orthogonality of complex waves in 2d.

2d waves

- Wavefronts point in the direction of ω
- A useful 2d-analog of the 1d angular frequency ω is the *norm* ||ω|| of the frequency vector ω:
 Set f(t) = ω/||ω|| t with real valued t, and you get a 1d wave in the direction of ω, where t plays the same role as in 1d.
- A projection parallel to a coordinate axis *i* yields a 1d phasor with angular frequency ω_i.

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Real waves are sums of complex waves

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 Using the relationships between cos, sin and the phasor,

$$\cos \varphi = \frac{1}{2} \cdot e^{i\varphi} + \frac{1}{2} \cdot e^{-i\varphi}$$

$$\sin \varphi = \frac{1}{2i} \cdot e^{i\varphi} - \frac{1}{2i} \cdot e^{-i\varphi} = -\frac{i}{2} \cdot e^{i\varphi} + \frac{i}{2} \cdot e^{-i\varphi},$$

we may write

►
$$\cos(\omega^T \mathbf{y}) = \frac{1}{2} \exp(i\omega^T \mathbf{y}) + \frac{1}{2} \exp(-i\omega^T \mathbf{y})$$

► $\sin(\omega^T \mathbf{y}) = -\frac{i}{2} \exp(i\omega^T \mathbf{y}) + \frac{i}{2} \exp(-i\omega^T \mathbf{y})$

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DFT on images

• Using the discrete phase vector $\boldsymbol{\omega}$ with components $\omega_1 = \frac{2\pi}{M} k$, $\omega_2 = \frac{2\pi}{N} \ell$ we get the

Discrete Fourier Transform (DFT) in 2d

$$S(k,\ell) = \sum_{m=0}^{M-1} \sum_{n=0}^{N-1} s(m,n) e^{-i2\pi \left(\frac{km}{M} + \frac{\ell n}{N}\right)}$$

Inverse Discrete Fourier Transform in 2d

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$$m{s}(m,n) = rac{1}{MN}\sum_{k=0}^{M-1}\sum_{\ell=0}^{N-1}m{S}(k,\ell)m{e}^{i2\pi\left(rac{km}{M}+rac{\ell n}{N}
ight)}$$

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DFT Symmetries

- 2. For real valued signals
 - in 1d: $S(k) = \bar{S}(T k)$
 - in 2d: $S(k, \ell) = \overline{S}(M k, N \ell)$
 - So the amplitude spectrum of real valued signals is symmetric.



DFT Symmetries

1. For discrete phasors

- ▶ 1d: The phasor p(k) is equal to the conjugate phasor $\bar{p}(T k)$.
- ▶ 2d: The phasor $p(k, \ell)$ is equal to the conjugate phasor $\bar{p}(M k, N \ell)$.
- So a "fast" phasor is the same as a "backwards spinning" slow phasor, and vice versa.
- ► Here, conjugate means elementwise.

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Practical implications of symmetries

- Because of this, the Fourier transform is usually presented with the origin in the center, and increasing (absolute) amplitudes away from the origin.
- All common software packages provide a function for this (such as "fftshift").



Spectrum example



 The amplitude spectrum is usually shown on a logarithmic scale, because it can span several orders of magnitude.

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More amplitude spectra (average over cross-sections on the right)

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from: Natural Image Statistics (Hyvarinen, Hurri, Hoyer; 2009)

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Spectrum example



 The amplitude spectrum is usually shown on a logarithmic scale, because it can span several orders of magnitude.

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Amplitude spectra of natural images

- Amplitude spectra of natural images show strong regularity.
- They typically are roughly of the form

 $\frac{1}{f}$

where *f* is the absolute value of frequency.

- Equivalently, the power spectrum $|S(f)|^2$ takes the form $\frac{1}{f^2}$
- This is known as **power law**.
- Power law relationships occur frequently in nature and they can be related to scale invariance.
- They can also be related to the covariance structure of natural images.



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Low-pass filtering in the frequency domain

An easy way to define a low-pass filter is to multiply low frequency components by zero, then to do the inverse transform to get back the resulting signal.

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However, this will yield ringing artefacts:



Fourier transform of a box

The box function is defined as

$$b_T(t) = \left\{ egin{array}{cc} 1 & ext{if} & -rac{T}{2} \leq t \leq rac{T}{2} \\ 0 & ext{else} \end{array}
ight.$$

Its Fourier transform is

$$B_{T}(f) = \int_{-\infty}^{\infty} b(t) e^{-i2\pi ft} dt$$
$$= \int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-i2\pi ft} dt \propto \frac{\sin(\pi Tf)}{\pi Tf} = \operatorname{sinc}(Tf)$$

 Similarly, the impulse response corresponding to a frequency "box" response is a sinc function, too.



Ringing



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Avoiding sharp transitions

One can relate the problem of ringing also to the presence of sharp edges in the filter.

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- Discontinuities in the signal require the presence of many spectral components (and vice versa).
- An informal rule of thumb for designing filters is therefore to avoid sharp edges.

Filterdesign

Filterdesign: "window method"

- 1. Specify desired ideal filter H_{id} , as well as *tolerances* in the frequency domain.
- 2. Compute the inverse Fouriertransform h_{id} of the desired filter. This will in general not be finite. Therefore:
- 3. Multiply h_{id} with a window *w*. Windows have been computed analytically and are tabulated wrt. to their tolerances in the frequency domain.
- 4. Implement the filter by convolving with mit $h_{\rm id} \cdot w$

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Gauss filter

$$g(x,y) = \frac{1}{2\pi\sigma} \exp\left(-\frac{1}{2\sigma^2}(x^2+y^2)\right)$$

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Gauss filter

The Fourier transform of a Gaussian is a Gaussian.

• The standard deviations are related by $\sigma_{\rm spectrum} = \frac{1}{\sigma}$

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 In practice: Use a sampled Gaussian or some other discrete approximation to the Gauss function.

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Gauss filter, $\sigma = 1.0$



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Example



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Gauss filter, $\sigma = 2.0$



